


Spectral Networks

[w/ D. Gaiotto, G. Moore]

Aim: to describe some new geometric objects — networks of curves drawn on a punctured Riemann surface C — and explain how they control the BPS spectrum of $\mathcal{N}=2$ theories $S[g=A_{K-1}, C, D]$

The result will look like s.t. that could come from string theory — but in fact it's obtained by purely field theory considerations (w/ optimistic assumptions)

Players:

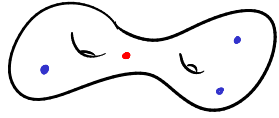
Riemann surface C 
 $K \geq 2$
 Marked points z_l $1 \leq l \leq n$
 Parameters $m_l^{(i)} \in \mathbb{C}$ $1 \leq l \leq n, 1 \leq i \leq K$
 (or more general "defects")

$\mathcal{N}=2$ SUSY QFT
 in $d=4$
 [Witten, GMN, Gaiotto]

Tuple $\vec{\phi} = (\phi_2, \dots, \phi_K)$
 ϕ_r a meromorphic r -differential on C
 with poles only at the z_l , residues det. by $m_l^{(i)}$

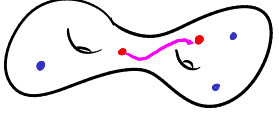
a point of the Coulomb branch

IR physics: governed by the curve $\Sigma = \{ (\lambda, z) \in T^*C : \lambda^K + \sum_{r=2}^K \lambda^{K-r} \phi_r = 0 \}$

Point $z \in C$ 

a surface defect

S_z

Path P from z_1 to z_2 
 Phase ϑ

a supersymmetric interface

S_{z_1}

S_{z_2}

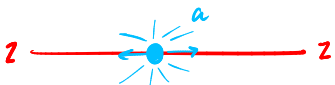
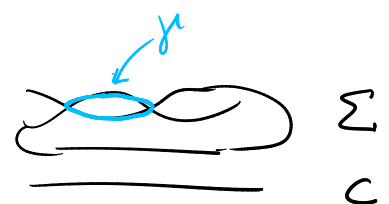
$L_{P, \vartheta}$

[AGTV, Gaiotto]

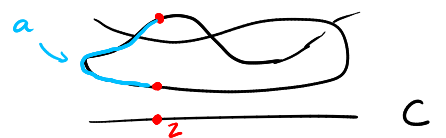
Now we may consider various sorts of BPS states. In particular:

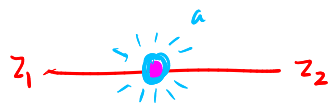


4d BPS states: charge $\gamma \in H_1(\Sigma, \mathbb{Z})$
 $Z = \oint_{\gamma} \lambda$
 count $\Omega(\gamma, \vec{\phi})$



2d BPS states: charge $a \in \mathbb{Z}$
 $Z = \int_{\gamma} \lambda$
 count $\mu(a, z, \vec{\phi})$

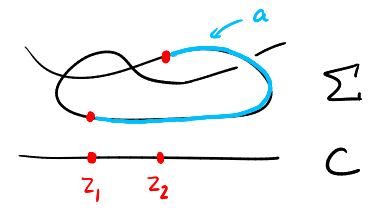




2d framed BPS states

$$Z = \int_a \lambda$$

count $\underline{\Omega}(a, P, \mathcal{D}, z, \bar{\Psi})$

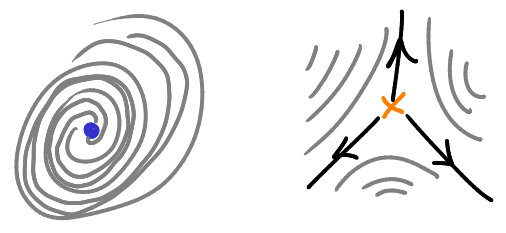


Now, how to compute the BPS spectra? By looking at their wall-crossing behavior + using just homotopy invariance, find that they are all governed by a single object:
spectral network $W(\mathcal{D}, \bar{\Psi})$ drawn on C .

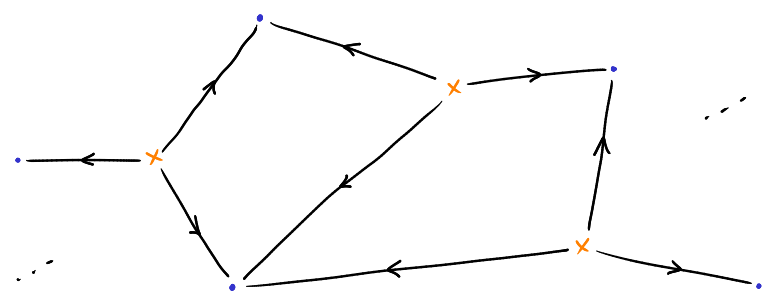
Begin with the special case $K=2$. So we have a curve C with a meromorphic quadratic differential ϕ_2 . Assume ϕ_2 has only simple zeroes.

Given any $\mathcal{D} \in \mathbb{R}/2\pi\mathbb{Z}$ define \mathcal{D} -trajectories of ϕ_2 to be paths on C along which $e^{-i\mathcal{D}} \sqrt{\phi_2}$ is real.

\mathcal{D} -trajectories make a singular foliation $F(\mathcal{D}, \phi_2)$ of C :
 $F(\mathcal{D}, \phi_2)$ is singular at poles and zeroes of ϕ_2 .



Focus on the 3 trajectories emanating from each zero; let $W(\mathcal{D}, \phi_2)$ be their union.

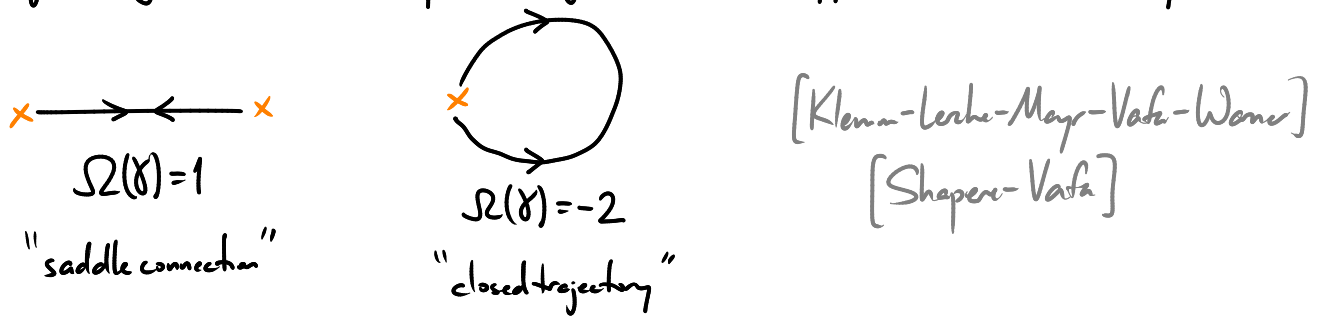


Then:

- The framed BPS counts $\bar{\Omega}(a, P, \mathcal{D}, \varphi_2)$ are determined by the intersections between P and $W(\mathcal{D}, \varphi_2)$. [Path lifting rule] (\Rightarrow revs of line ops on $\mathbb{R}^3 \times S^1$)
- When \mathcal{D} is varied, $W(\mathcal{D}, \varphi_2)$ sometimes jumps. \Rightarrow framed BPS states appear/disappear. How? By absorbing/emitting 4d BPS states! So:

$\Omega(\gamma, \varphi_2) \in \mathbb{Z}$ count the jumps of $W(\mathcal{D}, \varphi_2)$.

Concretely these jumps come from special trajectories which appear in $W(\mathcal{D}, \varphi_2)$ at special \mathcal{D} :



- An advantage of this point of view:
Varying along different paths in (\mathcal{D}, φ) parameter space, framed BPS states appear/disappear in different order. But the end result must be path independent!
 \Rightarrow Kontsevich-Soibelman wall-crossing formula for $\Omega(\gamma, \varphi_2)$.

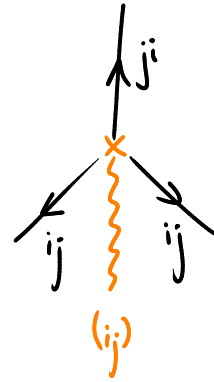
Now the general case, $K \geq 2$:

Locally, label the sheets of Σ by $i=1, \dots, K$

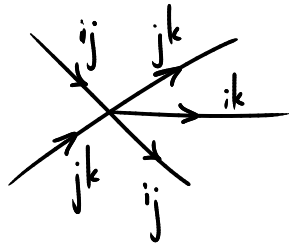
Then we get K locally defined hol 1-forms $\lambda^{(i)}$ on C .

Define an **ij -trajectory with phase ϑ** to be a real curve on C with $e^{-i\vartheta}(\lambda^{(i)} - \lambda^{(j)})$ real, positive.

Each branch point of type (ij) has 3 outgoing ij - or ji -trajectories

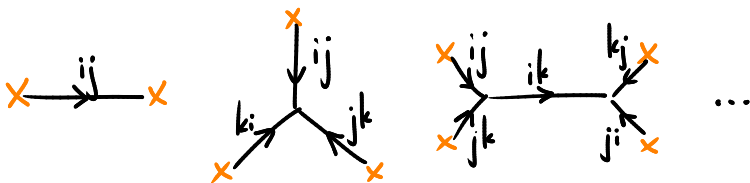


To define spectral network $W(\vartheta, \varphi)$: take all these trajectories plus extra ones generated at intersections:



As in $K=2$ case, this object determines all BPS counts!

In particular, its jumps as ϑ varies determine $\Omega(\gamma, \vec{\varphi})$. This is counting various kinds of "finite webs":



As before, these $\Omega(\gamma, \vec{\varphi})$ obey KSWCF as $\vec{\varphi}$ varies.

Views of line operators \Rightarrow generalizⁿ of FG coordinates assoc'd to triangulations